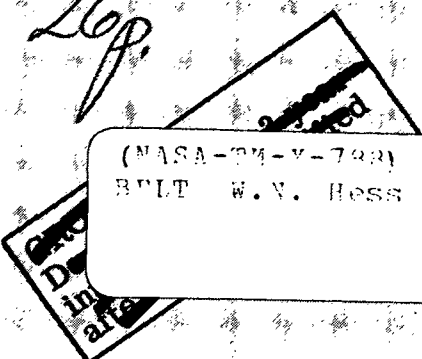


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# TECHNICAL MEMORANDUM

X-788

## THE ARTIFICIAL RADIATION BELT

W. N. Hess

Goddard Space Flight Center  
Greenbelt, Maryland

CLASSIFICATION CHANGED  
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NATIONAL AERONAUTICS AND SPACE ADMINISTRATION

WASHINGTON

April 1963

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CONFIDENTIAL - SECURITY INFORMATION

## THE ARTIFICIAL RADIATION BELT

by

Wilmot N. Hess

*Goddard Space Flight Center*

### SUMMARY

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The available information on the artificial radiation belt formed by the July 9, 1962, high altitude nuclear explosion is reviewed. Data from Injun (1961 02), Telstar I (1962 aε1), Traac (1961 aη2), and Ariel I (1962 o1) are combined to form one picture of the artificial belt. The data are consistent to about a factor of 3. The flux map obtained in this way is used to calculate the flux encountered by several satellites. These show reasonable agreement with data on solar cell damage. Preliminary data on particle lifetimes are presented. Particles at  $L > 1.30$  are expected to last several years on the basis of coulomb scattering. Crude calculations of shielding are made to indicate the doses received inside various vehicles.

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# THE ARTIFICIAL RADIATION BELT\*

by

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## INTRODUCTION

On July 9, 1962, at 0900:09 UT a nuclear explosion of about 1.4 megatons was carried out at 400 kilometers above Johnston Island in the Pacific Ocean. This explosion produced, as was expected, an artificial radiation belt. However, the intensities in this radiation belt are considerably higher than were expected. Three days after the explosion the U.S.-U.K. joint satellite Ariel I (1962 01) stopped transmitting. On August 2, Transit IV-B (1961  $\alpha\eta$ 1) stopped transmitting; Traac (1961  $\alpha\eta$ 2) stopped on August 14. Instruments on Ariel I, Traac, and Injun (1961 02) showed large particle fluxes shortly after the explosion. It took about a month to start getting some grasp of the characteristics of the new radiation belt. This is a status report on the new belt as of September 12.

## AVAILABLE DATA

The information that is available to form a picture of the new radiation belt comes mostly from particle detectors on the Ariel I, Injun, Traac, and Telstar I (1962  $\alpha\epsilon$ 1) satellites. In addition to these data we can use the observed solar cell damage on satellites as an integral measurement of the trapped electron flux. Also, some data are available from dosimetry measurements.

Some of the original data about the enhanced trapped particle fluxes after the July 9 explosion came from the x-ray detector on the Ariel I satellite (private communication from A. Willmore, University College, London). This instrument was not designed to count charged particles and therefore its efficiency is uncertain. The data from it are quite useful in studying the time decay of the trapped flux and in locating contours of constant flux in B-L space.

Data received by the shielded 213 GM counter on Injun have been analyzed to give the first picture of the new radiation belt (Reference 1). This counter is the background channel of the magnetic spectrometer, SpB. It has 3-1/2 gm/cm<sup>2</sup> of Pb shielding and about 1 gm/cm<sup>2</sup> of wall and miscellaneous shielding. It was supposed to give the penetrating background to be subtracted from the other channels

\*Title unclassified. An abridged version of this report will appear in the *Journal of Geophysical Research* and will be published as NASA Technical Note D-1687 under the title "The Artificial Radiation Belt Made on July 9, 1962."

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of the spectrometer. This detector is now called on to provide quantitative information, and it has been calibrated after the fact. It is nearly omnidirectional. Fluxes are obtained from the count rates by dividing by  $G_0 = 0.11 \text{ cm}^2$ . Other detectors on Injun also give useful data sometimes, but often they are saturated and not usable. So far, little data have been analyzed from any Injun detectors except SpB.

Telstar I has on it a solid state p-n junction detector with pulse height analysis that selects electrons in different energy ranges from 0.2 to 1 Mev (private communication from W. Brown, Bell Telephone Laboratories). A lot of data have been reduced from Telstar I for two channels of the electron detector. This detector has given all the data currently available at high altitudes. It is directional, with an aperture half-angle of about 10 degrees. The fluxes are made omnidirectional by multiplying by the appropriate solid angle factor and then using a factor between 1 and 2 to correct, roughly, for the nonisotropic angular distribution.

Traac has a 302 GM counter shielded by  $0.265 \text{ gm/cm}^2$  of Mg, which will count electrons of energy above 1.5 Mev (private communication from G. Pieper and L. Frank, Applied Physics Laboratory). It is essentially omnidirectional. Fluxes are obtained by dividing by  $G_0 = 0.75 \text{ cm}^2$  and correcting for saturation for high count rates.

## ANALYSIS OF THE DATA

The data from these four satellites must be combined to form one overall picture of the artificial radiation belt. To do this assume that the energy spectrum of the electrons being counted is a fission spectrum. This is certainly the best guess. We will compare the data on this basis and see if there is agreement in the regions where direct comparison is possible. The fission energy spectrum  $N(E)$  is shown in Figure 1, curve A. A calibration of the Telstar I detectors at the Los Alamos Scientific Laboratory in a fission electron beam gives  $f$ , the fraction of fission electrons counted by the detectors, equal to  $1/2.8$  for the 240-340 kev channel and  $1/6.0$  for the 440-680 kev channel.

For Injun we have the experimentally determined factor  $1/f$  of several thousand, by comparison of two detectors on board. The 213 GM counter has also been calibrated at Los Alamos with a fission electron spectrum (private communication from A. Petschek, H. Motz, and R. Taschek, Los Alamos Scientific Laboratory), and the factor  $f$  determined this way is  $1/4000$ . We will use this factor in the present analysis. The Los Alamos tests show that the detector counts bremsstrahlung from electrons of several Mev rather than direct penetrating electrons. (If the shield had been carbon rather than lead, the counter would have counted direct penetrating electrons.)

For Traac,  $f$  is determined by considering the penetration of electrons through the detector shield of  $0.265 \text{ gm/cm}^2$  of Mg and through the wall of  $0.400 \text{ gm/cm}^2$  of stainless steel. Using the range straggling data (Reference 2) for Al we can get the fraction of electrons that penetrate a shield of given thickness, as shown in Figure 2. The expression for the extrapolated range  $R$  is

$$R = 0.526 E - 0.094 .$$

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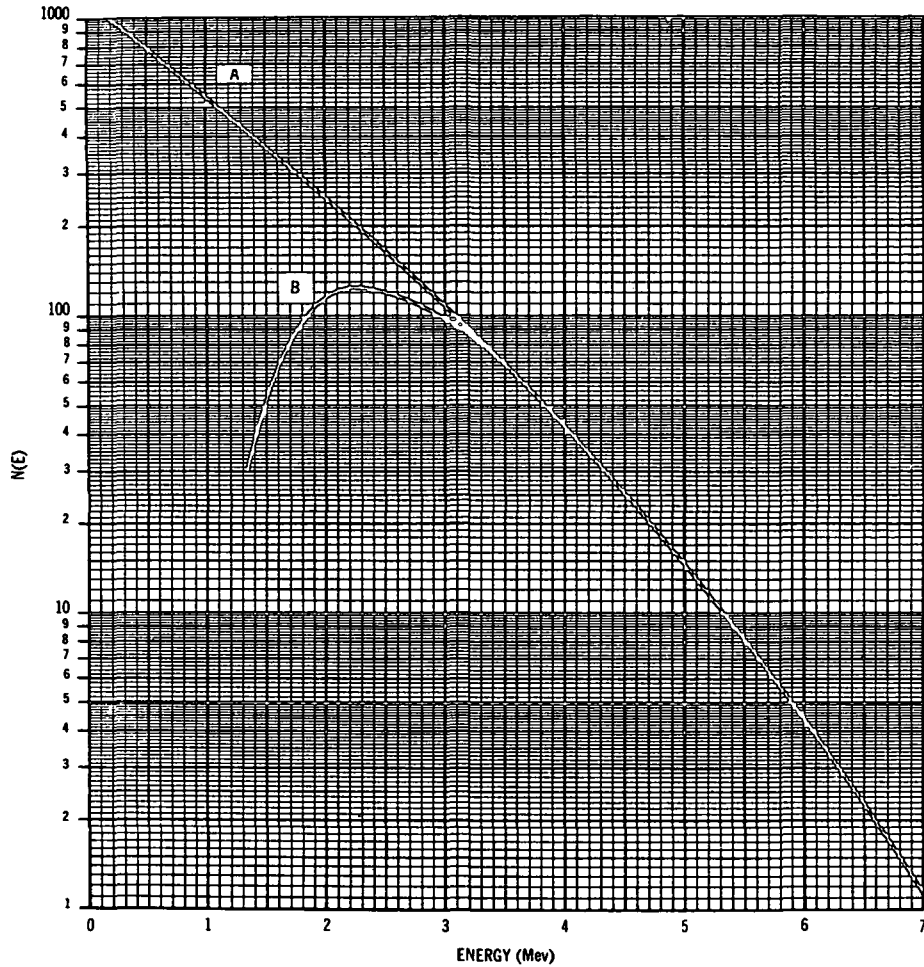


Figure 1—Curve A is the fission energy spectrum and curve B the transmission energy spectrum for the Traac GM counter (0.66 gm/cm<sup>2</sup> wall).

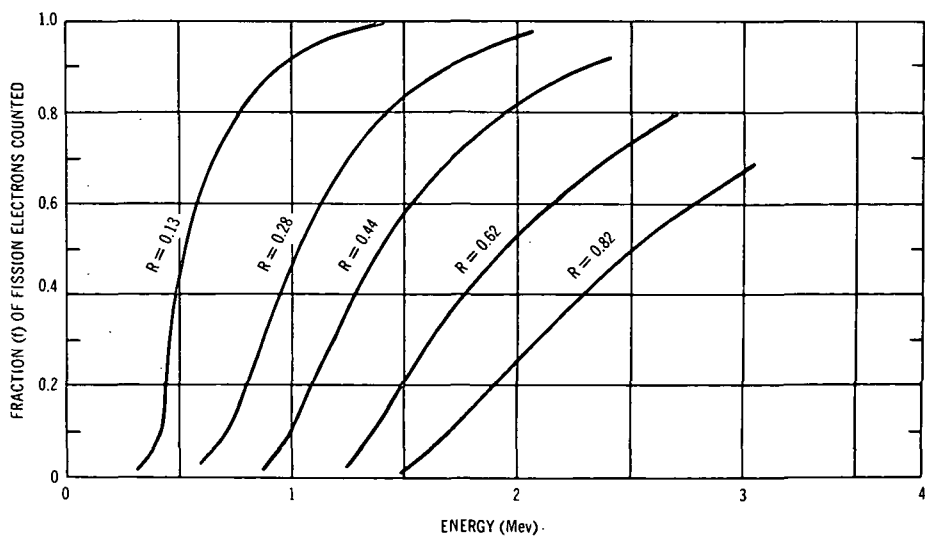


Figure 2—The fraction of electrons of different energies that penetrate different shield thicknesses of Al.

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This yields the absorber thickness that gives 10 percent transmission for electrons of energy E. For 50 percent transmission we multiply the energy by 1.38, and for 80 percent transmission increase the energy by a factor of 1.92. In this way we get the electron transmission spectrum, curve B in Figure 1. The energies of the transmitted electrons are different from curve B, but the number transmitted is given correctly. The integral under this curve gives  $f = 1/5.5$  for the Traac counter. More information on shielding calculations is given in Appendix A.

Using the factors for the several detectors, we can calculate the total flux of fission electrons. In order to compare the different detectors, the total flux along several field lines (actually narrow ranges of L) has been plotted for different values of B (Figure 3). These plots show that the different detectors agree fairly well in flux values. Avoiding the first day after the nuclear explosion (labeled by the number 0 inside the symbols on the graphs) we can see quite smooth trends in the data. The flux from Telstar I may be as much as twice as high as Injun fluxes. Traac and Injun agree quite well where comparisons are possible. In general, the data shows agreement to a factor of 2.

This agreement of the data shows two things: First, because the detectors give internally consistent results it seems likely that all the detectors are giving accurate information. Secondly, the assumption that the electrons have a fission energy spectrum appears to be correct. Of course it is possible that the energy spectrum is not a fission spectrum and also that the detectors are not in agreement, but it would have to be a peculiar combination of such effects that would give the agreement shown here. A comparison of the four channels of the Telstar I electron detector also indicates that the energy spectrum is fission-like up past 1 Mev.

## FLUX PLOTS

Now that it has been demonstrated that the energy spectrum is essentially a fission spectrum at least in the region of data overlap we can use all the counter data to construct a composite flux map in B-L space. As McIlwain has shown, these magnetic coordinates are the best way of organizing data about trapped particles (Reference 3). L is constant along a field line in space and, for a dipole, is the distance from the center of the earth to the equatorial crossing of the line, in units of earth radii (Figure 4). Values of L are calculated from the real values of the earth's field.

In constructing the flux map for  $B > 0.15$  gauss and for  $L < 2.0$  earth radii the graphs in Figure 3 are used to locate the flux contours. The experimental data outside this B-L region are essentially all from Telstar I. There are several weeks data from Telstar I and considerable redundancy. The map made this way is quite complete. The data available in early September gave the flux map in Figure 5. This map is for about 1 week after the explosion. There was considerably more flux at low altitudes at early times.

This same data plotted in R- $\lambda$  coordinates, where

$$B = \frac{M}{R^3} \sqrt{4 - \frac{3R}{L}},$$
$$R = L \cos^2 \lambda,$$

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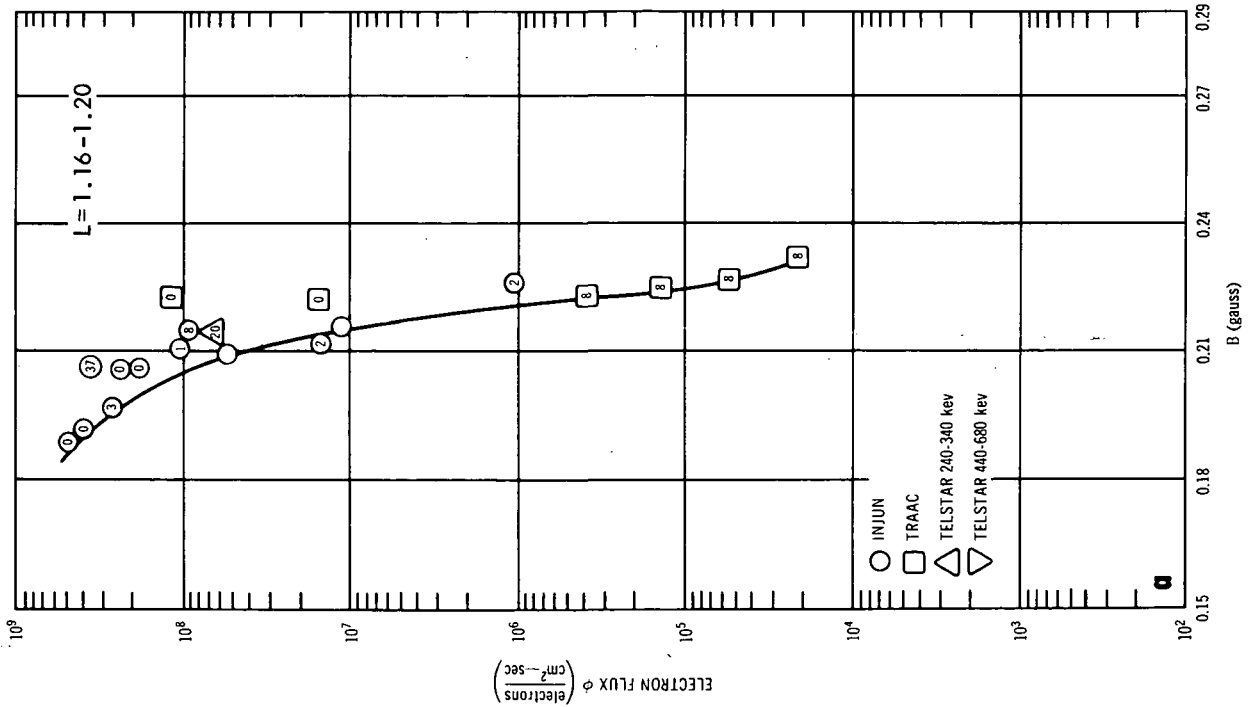
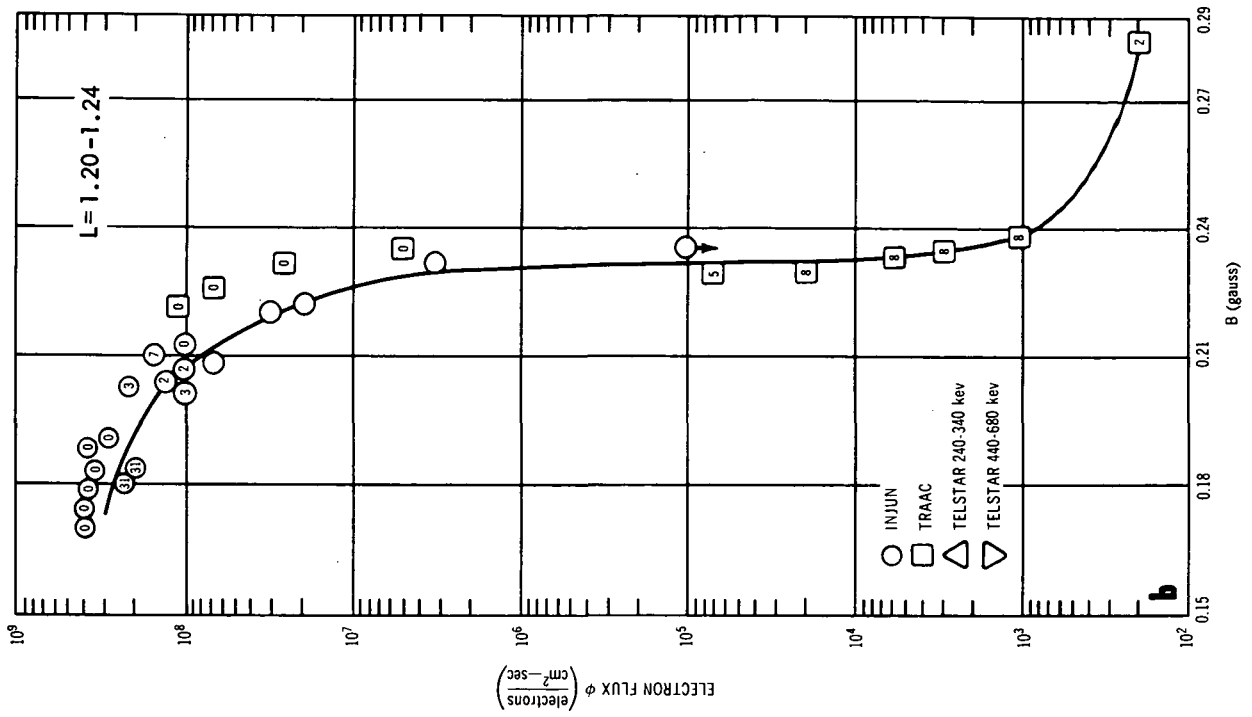


Figure 3—The electron flux distributions along different field lines.



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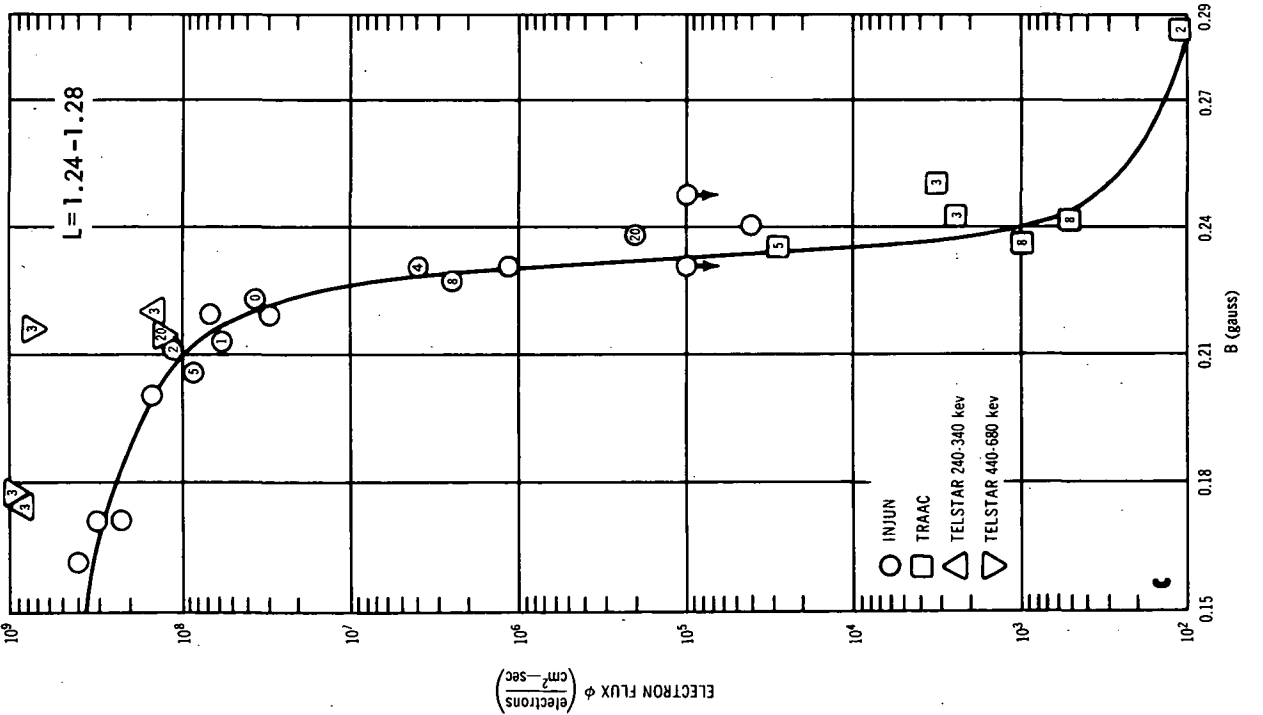
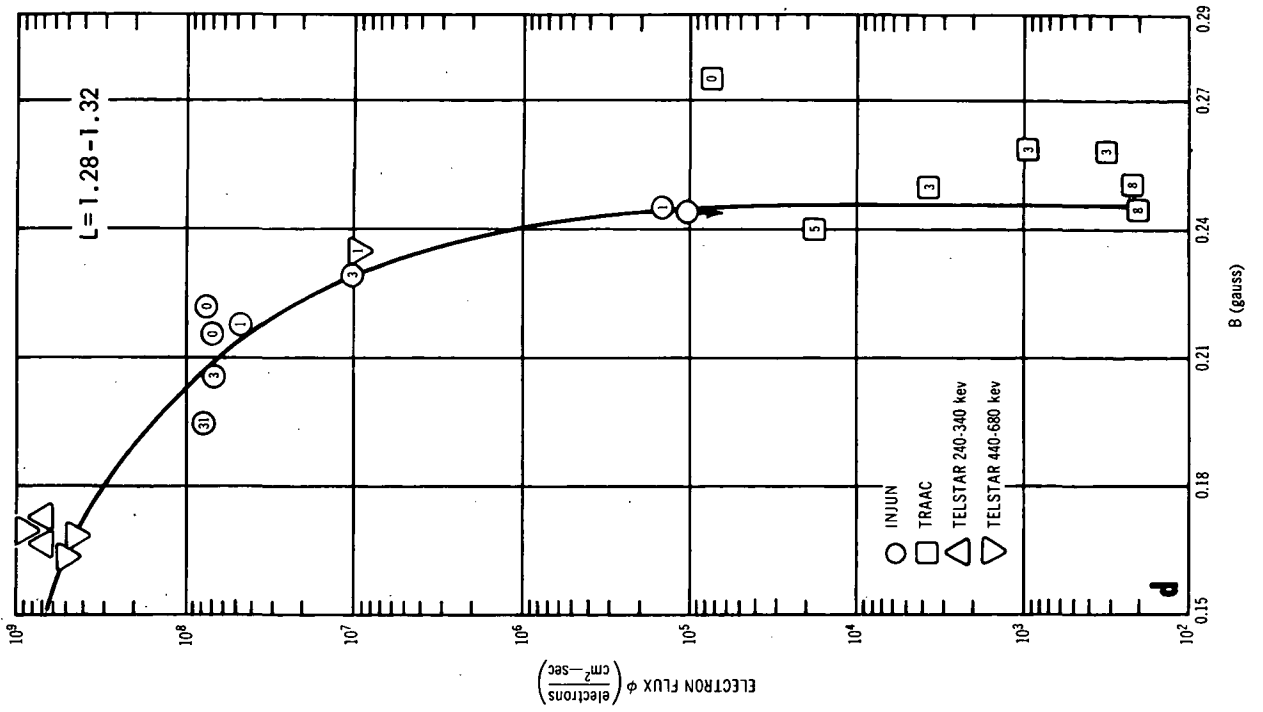


Figure 3 (continued)—The electron flux distributions along different field lines.

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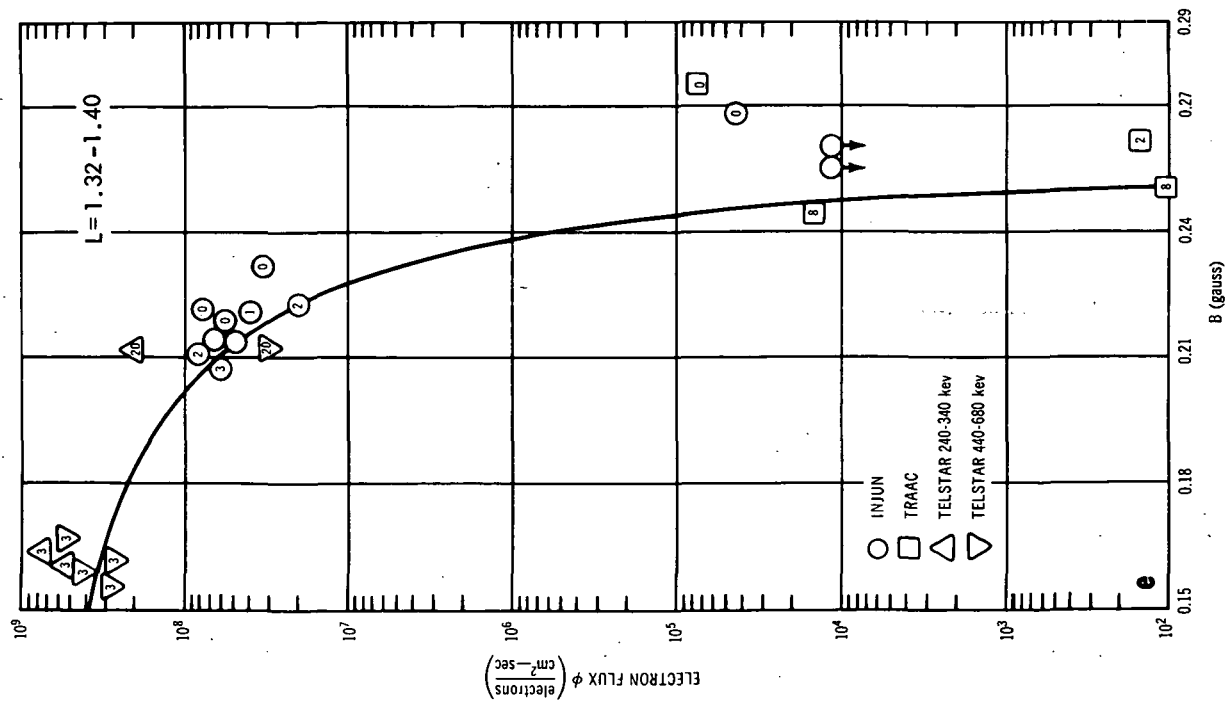
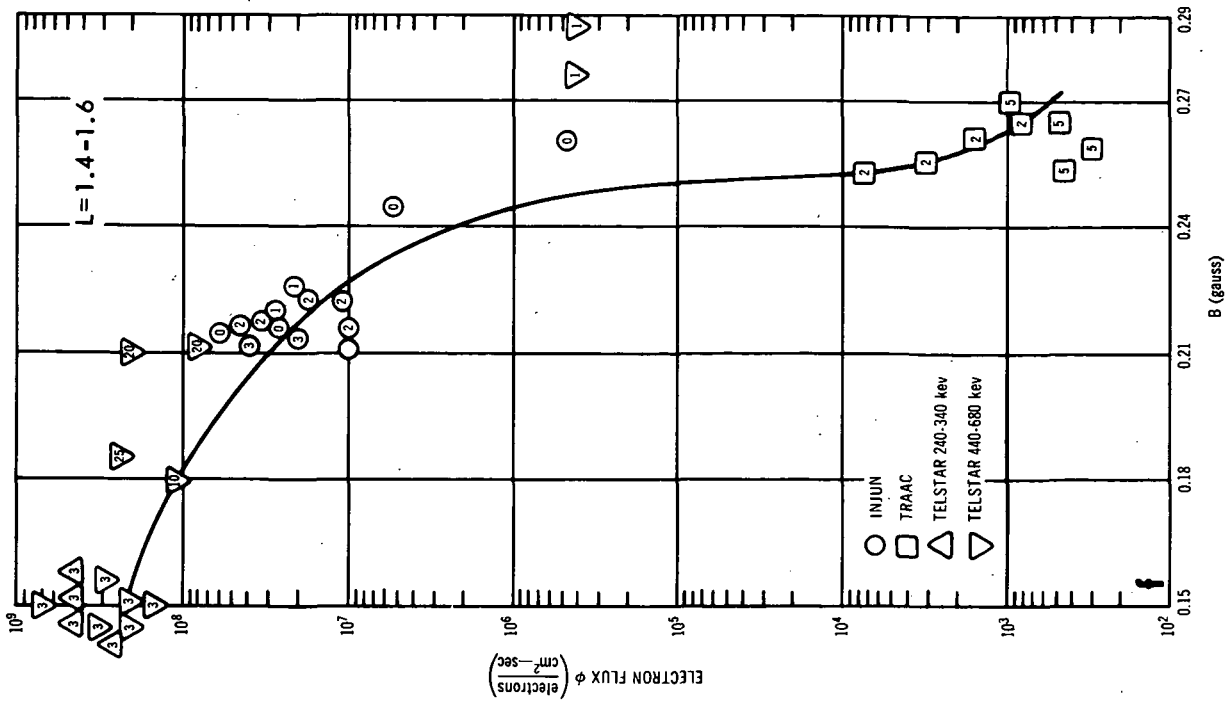


Figure 3 (continued)—The electron flux distributions along different field lines.

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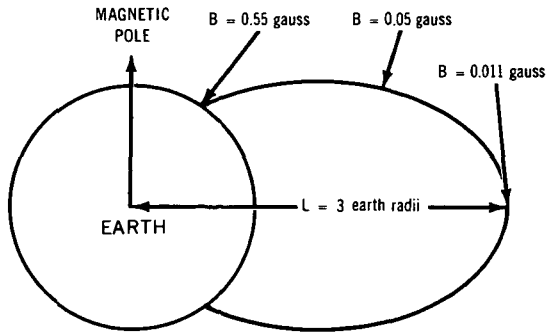


Figure 4—The B-L magnetic coordinate system.

gives an equivalent dipole representation of the earth's field (Figure 6). The maximum electron flux is about  $2 \times 10^9$  elec/cm<sup>2</sup>-sec. Integrating to get the total number of electrons stored in the field we find

$$\int \phi dV = 2 \times 10^{26} \text{ electrons.}$$

About 60 percent of these electrons lie inside the  $3 \times 10^8$  contour. It is not certain what fraction of these electrons are bomb-induced and what fraction are natural electrons. In this region around  $L < 1.5$  the energy spectrum seems softer than a fission spectrum.

The B-L flux map when plotted in terms of geographic coordinates gives the flux contours for different altitudes shown in Figure 7.

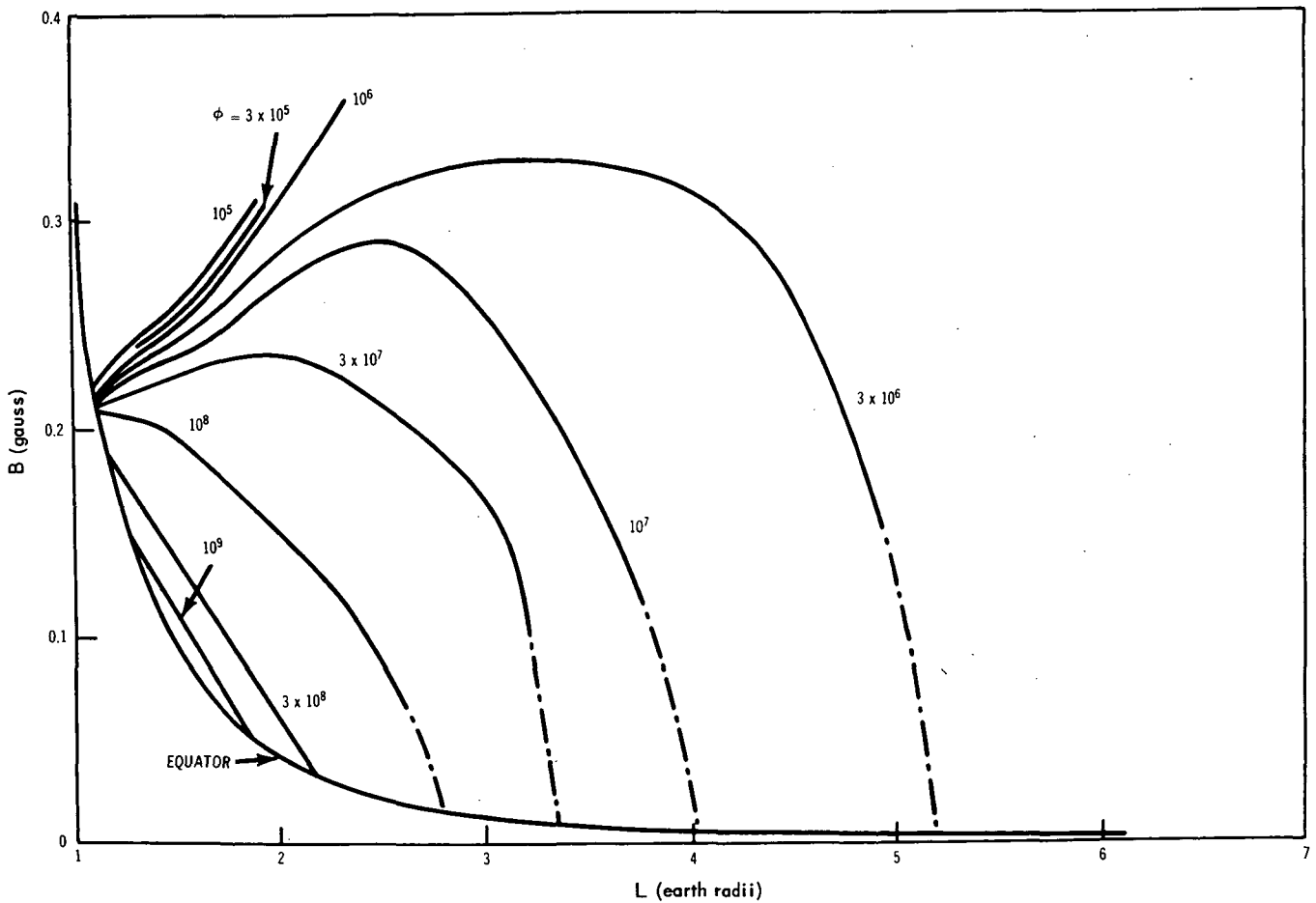


Figure 5—The B-L map of electron fluxes.

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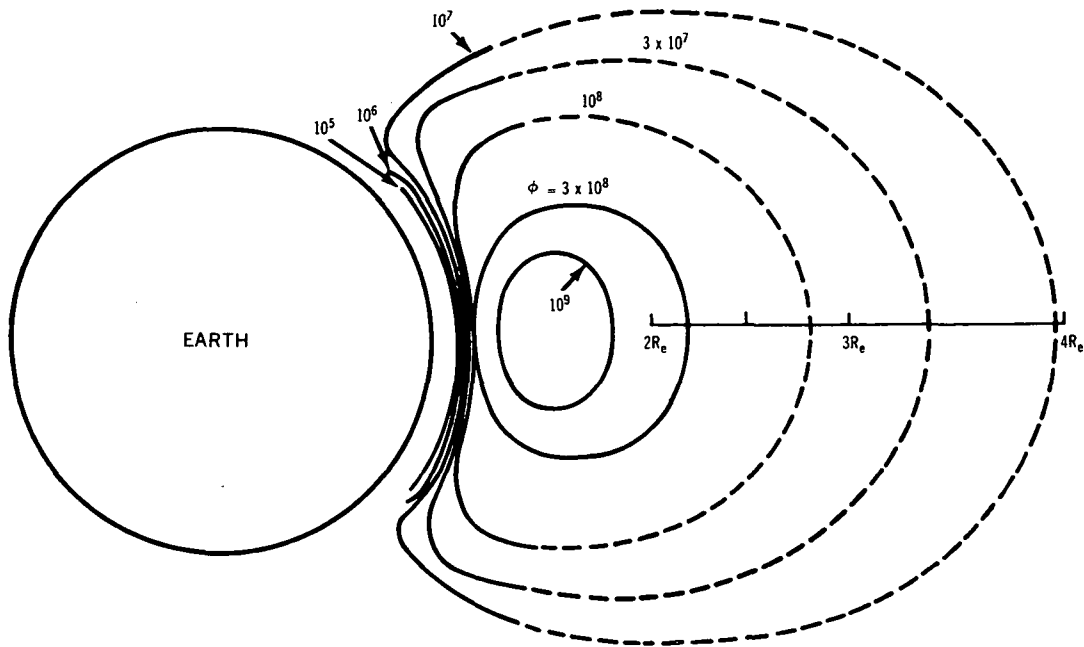


Figure 6—The R-λ map of electron fluxes (an ideal dipole representation of the earth's field).

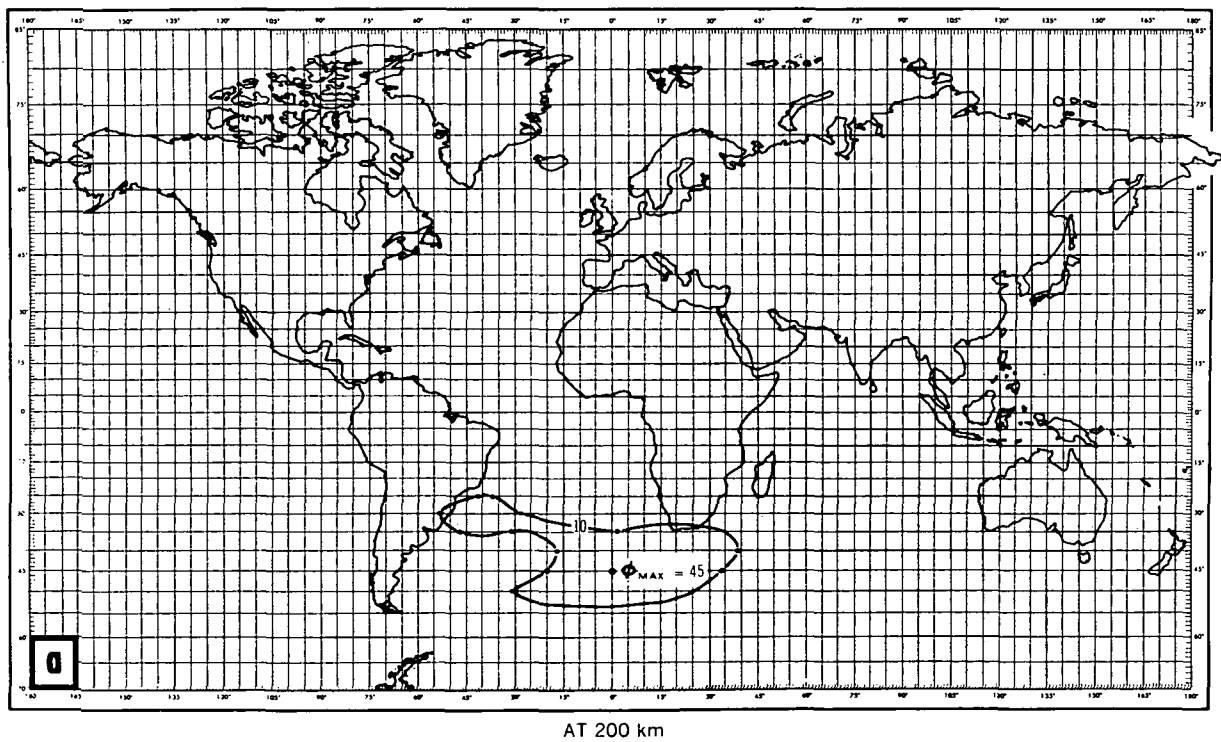
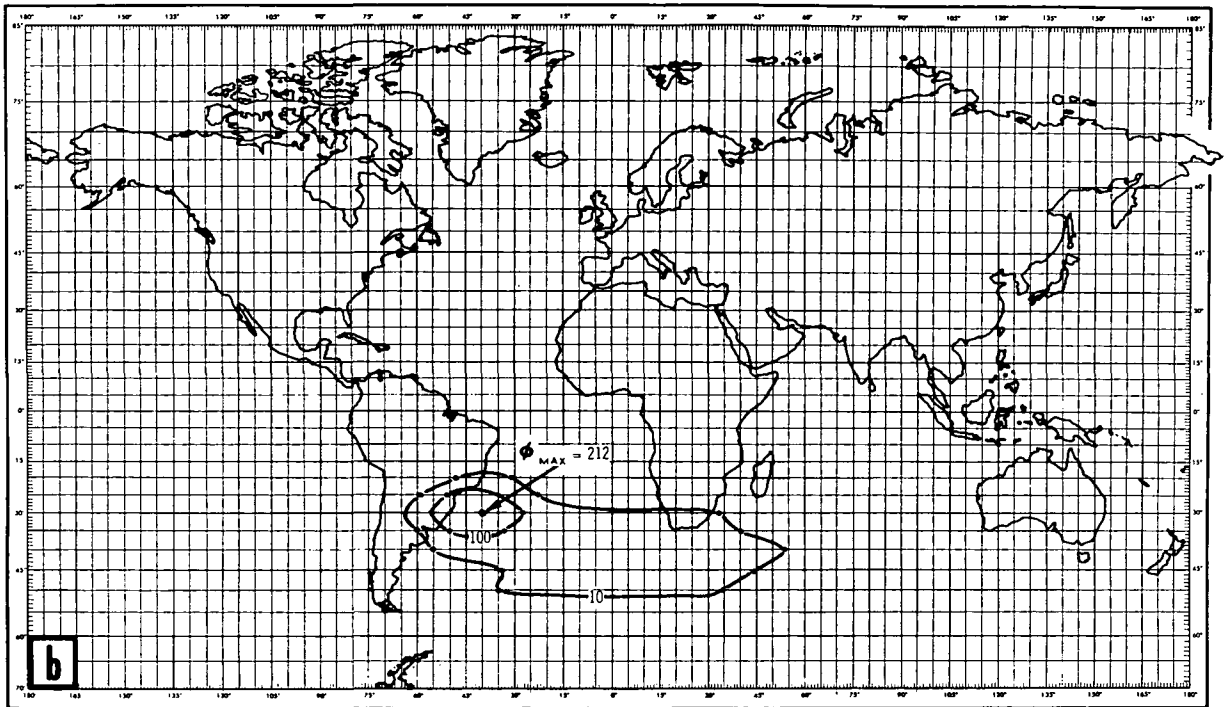
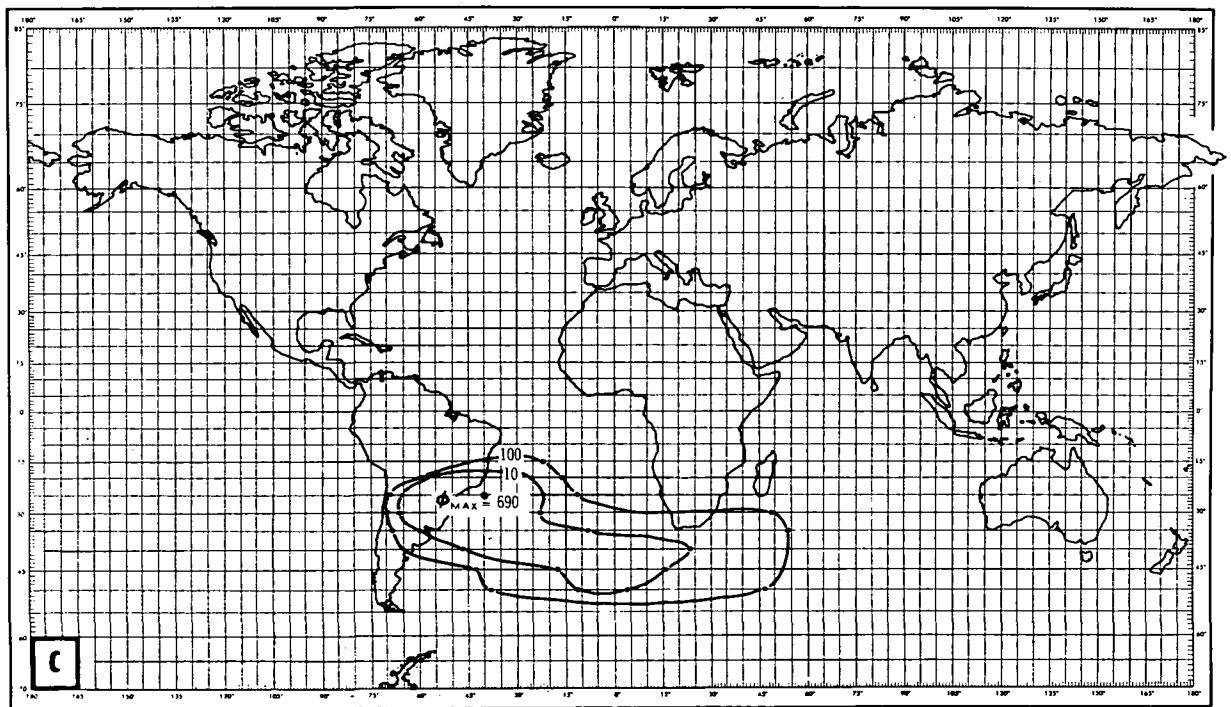


Figure 7—Electron flux maps at different altitudes above the earth's surface. Flux is in units of  $10^5$  electrons/cm<sup>2</sup>-sec.



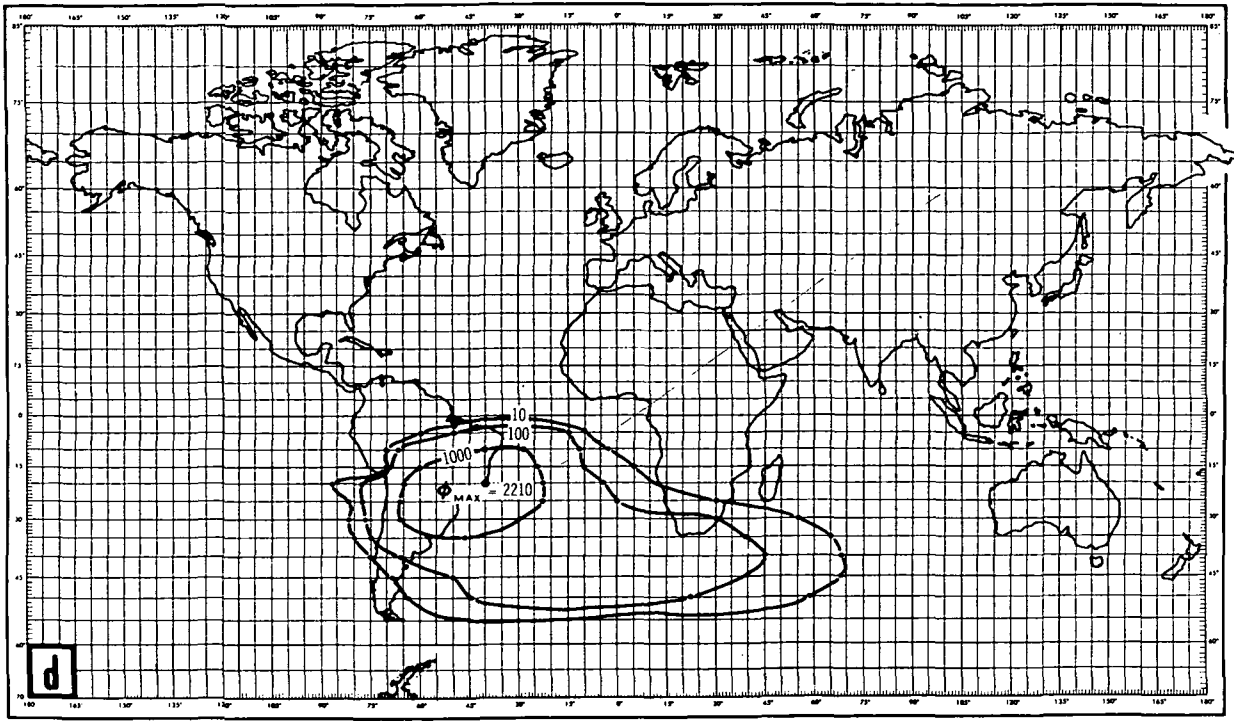
AT 300 km



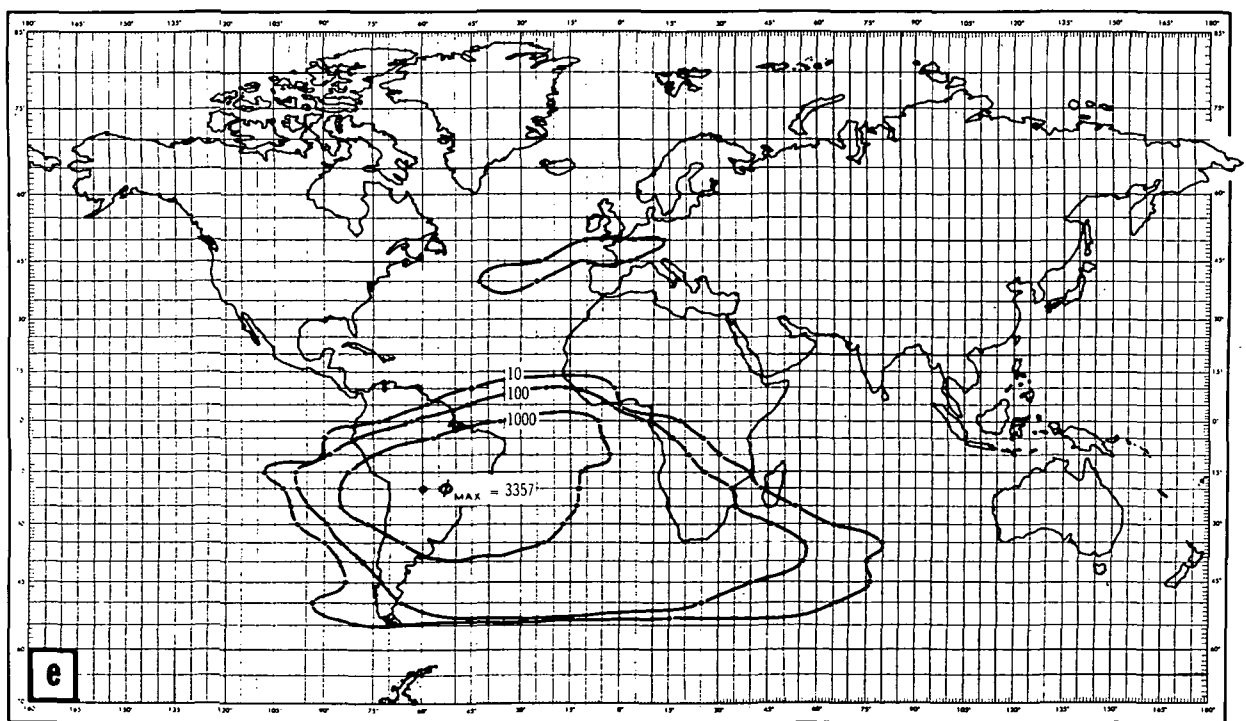
AT 400 km

Figure 7 (continued)—Electron flux maps at different altitudes above the earth's surface. Flux is in units of  $10^5$  electrons/cm<sup>2</sup>-sec.

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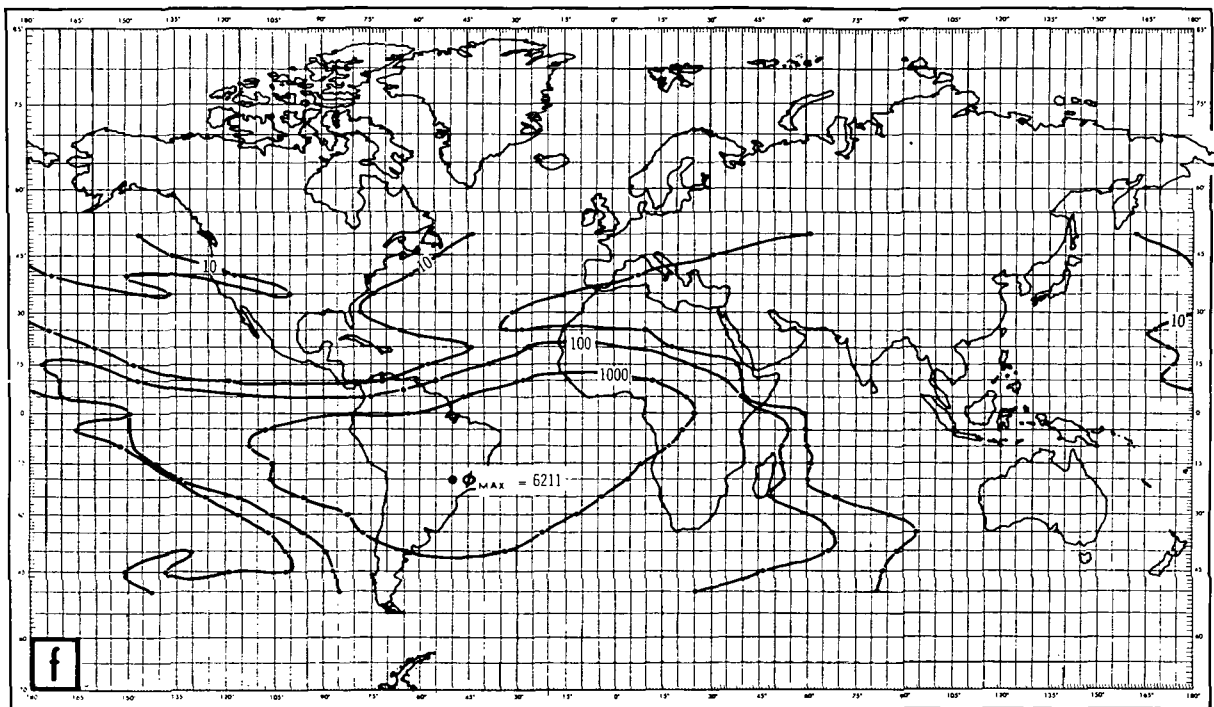
AT 600 km



AT 800 km

Figure 7 (continued)—Electron flux maps at different altitudes above the earth's surface. Flux is in units of  $10^5$  electrons/cm<sup>2</sup>-sec.

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AT 1000 km

Figure 7 (continued)—Electron flux maps at different altitudes above the earth's surface. Flux is in units of  $10^5$  electrons/cm<sup>2</sup>-sec.

## VEHICLE-ENCOUNTERED FLUXES

A machine code has been developed which calculates the total number of electrons/cm<sup>2</sup> in the artificial radiation belt that strike a vehicle in space. This is done by calculating a point on the vehicle trajectory, transforming to B-L coordinates, looking up the electron flux, and integrating along the vehicle orbit. This has been performed and the encountered fluxes have been determined for all of the vehicles listed in Tables 1 and 2. These fluxes have been transformed into r/day by using  $2.5 \times 10^7$  electrons/cm<sup>2</sup> = 1r. The orbital elements of these vehicle trajectories are given in Table 3.

From the encountered fluxes in Table 1, we can learn several things. Let us first consider solar cell damage. The Bell Telephone Laboratories staff have studied this problem in considerable detail and prepared Figure 8, which shows how different type cells are damaged by 1 Mev electron irradiation (Reference 4 and private communication from the authors of that paper). Above about 0.5 Mev the electron damage is essentially independent of energy. Some care must be exercised in using this chart because of the variation in the characteristics of solar cells. We will assume all the electrons in the flux spectrum in Figure 5 are greater than 0.5 Mev in estimating the solar cell damage.

About 20 percent degradation was needed by the blue sensitive p-on-n type cells on Ariel I to produce the observed power supply damage (private communication from A. Franta, Goddard Space

Table 1

Calculations on Fluxes Encountered by Satellites Moving Through the Artificial Radiation Belt.

Datum	Ariel I	Traac and Transit IV-B	Telstar I	Tiros V	Orbiting Solar Observatory I (1962 ζ1)	Relay*	Polar Orbiting Geophysical Observatory (Pogo)*
Perigee (km)	390	960	952	590	552	1343	257
Apogee (km)	1210	1106	5660	971	594	5555	931
Inclination (degrees)	54	32	45	58	33	50	90
Altitude (km) at 30° S Lat 30° W Long	1067	1000	5138 1758	963	594	4371 1516	804
Calculated r/day Outside Vehicle	110,000	180,000	800,000	46,000	27,000	$1.1 \times 10^6$	22,000
Code Output from Machine	$\frac{3.1 \times 10^9}{4}$	$\frac{5.0 \times 10^9}{4}$	$\frac{2.2 \times 10^{10}}{4}$	$\frac{1.28 \times 10^9}{4}$	$\frac{7.5 \times 10^8}{4}$	$\frac{3.0 \times 10^{10}}{4}$	$\frac{6.2 \times 10^8}{4}$
Length of Machine Run in Satellite Days	4	4	4	4	4	4	4
Electrons/cm <sup>2</sup> - day	$2.8 \times 10^{12}$	$4.5 \times 10^{12}$	$2.0 \times 10^{13}$	$1.15 \times 10^{12}$	$6.8 \times 10^{11}$	$2.7 \times 10^{13}$	$5.6 \times 10^{11}$
Protons/cm <sup>2</sup> - day	—	—	—	—	—	—	—

\*To be launched.



Table 1 (continued)  
 Calculations on Fluxes Encountered by Satellites Moving Through the Artificial Radiation Belt.

Datum	1000 km Polar Orbit†	800 km Polar Orbit	Orbiting Astronomical Observatory*	SERB*	MA-7 (1962 $\tau$ 1)	Vostok III (1962 $\mu$ 1)	Apollo* (round trip)
Perigee (km)	1000	800	802	278	160	159	—
Apogee (km)	1000	800	817	16,668	264	223	—
Inclination (degrees)	90	90	31	17.0	33	64.94	—
Altitude (km) at 30° S Lat 30° W Long	1000	755	810	—	261	—	—
Calculated r/day Outside Vehicle	80,000	27,000	—	—	100†	—	—
Code Output from Machine	$\frac{2.2 \times 10^9}{4}$	$\frac{1.04 \times 10^9}{4}$	$\frac{2.2 \times 10^9}{4}$	$\frac{1.4 \times 10^{10}}{4}$	$3.45 \times 10^6$	$\frac{8.6 \times 10^6}{4}$	—
Length of Machine Run in Satellite Days	4	4	4	4	0.4 (9.5 hours)	4	—
Electrons/cm <sup>2</sup> - day	$2 \times 10^{12}$	$9 \times 10^{11}$	$2 \times 10^{12}$	$1.2 \times 10^{13}$	$0.24 \times 10^{10}†$	$8 \times 10^9$	$6.3 \times 10^{12}$
Protons/cm <sup>2</sup> - day	—	—	—	—	—	—	—

\*To be launched.

†Similar to Nimbus, the Fixed Frequency Topside Sounder, the Swept Frequency Topside Sounder, and the Polar Ionosphere Beacon (which have been or will be launched).

‡Six orbits only.

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Table 2

Calculations on Fluxes Encountered by Department of Defense Satellites Moving Through the Area Occupied by the Artificial Radiation Belt Before and After Its Creation.

Datum	1129 23 June	1151 28 June	1130 21 July	1131 28 July	2 August	1133	2000 Nautical Mile Polar Orbit
Perigee (km)	205.558	200	204	204	204	205.0	3704
Apogee (km)	314.002	713	393	420	423	672.8	3704
Inclination (degrees)	75	76	70	70	82	82	90
Altitude (km) at 30° S Lat 30° W Long	221	250	384	412	415	—	2000
Calculated r/day Outside Vehicle	0.015	0.25	1800	2400	2400	—	—
Measured r/day Extrapolated to Outside of Vehicle	0.02	0.30	~2000	~2000	~2000	—	—
Code Output from Machine	$\frac{116.67}{4}$	$\frac{3584}{4}$	$\frac{4.87 \times 10^7}{4}$	$\frac{6.6 \times 10^7}{4}$	$\frac{6 \times 10^7}{4}$	$\frac{3 \times 10^8}{4}$	$\frac{8.5 \times 10^5}{4}$
Length of Machine Run in Satellite Days	4	4	4	4	4	4	4
Electrons/cm <sup>2</sup> - day	—	—	$4.5 \times 10^{10}$	$6.0 \times 10^{10}$	$5.4 \times 10^{10}$	$2.8 \times 10^{11}$	$2.05 \times 10^{13}$
Protons/cm <sup>2</sup> - day	$1.05 \times 10^5$	$3.2 \times 10^6$	—	—	—	—	$7.62 \times 10^8$

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[REDACTED]

Table 3  
Orbital Elements of Various Space Vehicles.

Element	Ariel I	Traac and Transit IV-B	Telstar I	Tiros V	OSO I	Relay
Epoch (days, hours, min, sec)	190, 9, 0, 0	190, 4, 3, 46.506	191, 8, 51, 0	190, 9, 0, 0	190, 9, 0, 0	305, 0, 0, 0
Semimajor Axis (earth radii)	1.1254	1.1618	1.5182	1.1224	1.0900	1.5407
Eccentricity	0.05714	0.009922	0.2430	0.02663	0.003012	0.2143
Inclination (degrees)	53.866	32.423	44.803	58.102	32.855	50.0003
Right Ascension of Ascending Node (degrees)	-24.881	96.434	-156.222	-75.536	154.502	163.708
Argument of Perigee (degrees)	-9.2537	-51.6890	164.811	118.014	139.136	-167.526
Mean Anomaly (degrees)	-86.8833	0.0001	1.1684	-194.11968	-164.5453	7.8219

Table 3 (continued)

Element	Pogo	1000 km Polar Orbit	Vostok III	Vostok IV	MA-7	200 km Circular Mercury Orbit	1129 23 June
Epoch (days, hours, min, sec)	82, 3, 55, 32.101	190, 9, 0, 0	223, 9, 50, 12.768	226, 1, 7, 14.016	268, 14, 0, 0	82, 3, 55, 32.101	174, 0, 37, 58
Semimajor Axis (earth radii)	1.0931	1.1568	1.0299	1.02922	1.0331	1.0314	1.0407295
Eccentricity	0.04830	$0.1490 \times 10^{-7}$	0.004890	0.00357	0.008552	$0.3057 \times 10^{-6}$	0.008169
Inclination (degrees)	90.001	90.000	64.940	64.99	32.546	33.000	75.099
Right Ascension of Ascending Node (degrees)	-73.806	-158.175	122.446	110.231	75.069	54.328	-11.0165
Argument of Perigee (degrees)	-19.408	180.000	81.049	89.496	78.188	-72.720	140.952
Mean Anomaly (degrees)	2.1956	0.0000	-80.5031	-89.087	7.6908	-177.6883	14.9695

[REDACTED]

Table 3 (continued)

Orbital Elements of Various Space Vehicles.

Element	1151 28 June	344	345	360	1132	698	1153
Epoch (days, hours, min, sec)	179, 1, 16, 22.52	202, 1, 3, 51.200	209, 1, 37, 20.800	214, 1, 24, 41.960	234, 1, 30, 0	233, 18, 19, 22.300	234, 1, 30, 0
Semimajor Axis (earth radii)	1.0715	1.0467	1.0489	1.0491	1.0721	1.1151	1.0513
Eccentricity	0.03754	0.01414	0.01619	0.01635	0.02327	0.003290	0.01818
Inclination (degrees)	76.058	70.297	71.085	82.251	81.789	98.410	65.017
Right Ascension of Ascending Node (degrees)	2.9295	26.456	41.296	35.380	56.667	-77.020	68.583
Argument of Perigee (degrees)	142.022	152.339	149.498	145.899	150.384	164.782	147.175
Mean Anomaly (degrees)	13.6621	2.9008	5.8434	11.0481	6.2969	-160.2016	6.5474

Flight Center). This would be caused by about  $10^{13}$  electrons/cm<sup>2</sup> according to Figure 8. About seven days after the nuclear explosion, this flux would have been achieved (Table 1 gives  $2.8 \times 10^{12}$  electrons/cm<sup>2</sup>-day for Ariel I, of which half hit the face of the cells). The Ariel I power supply started malfunctioning in 3-1/2 days. This is quite good agreement.

Traac and Transit IV-B also had blue sensitive p-on-n solar cells, but it would take  $3 \times 10^{14}$  electrons/cm<sup>2</sup> to cause malfunction, because the cells were lower efficiency cells (private communication from R. Fischell, Applied Physics Laboratory). Table 1 gives  $4.5 \times 10^{12}$  electrons/cm<sup>2</sup> encountered per day. Half of these electrons hit the face of the cells. Traac stopped transmitting in 36 days and Transit IV-B in 24 days. Using 30 days as the average, we get a total encountered flux of  $0.7 \times 10^{14}$  electrons/cm<sup>2</sup>, in moderate agreement with that required to produce damage.

Telstar I used the much more damage resistant n-on-p cells, because it was to routinely fly through the inner radiation belt protons. Even with the artificial radiation belt, its power supply lifetime is expected to be considerably longer than 1 year.

The Telstar I solar cells are degrading at a rate that would be produced by  $6 \times 10^{12}$  electrons/cm<sup>2</sup>-day of 1 Mev hitting the bare cells (private communication from W. Brown, Bell Telephone Laboratories). This corresponds to about  $1.8 \times 10^{13}$  electrons/cm<sup>2</sup>-day incident on the outside of the 30 mil sapphire covers. Our calculations give  $1/2 \times 2 \times 10^{13} = 1 \times 10^{13}$  electrons/cm<sup>2</sup>-day hitting the cells. The observed solar cell degradation on Telstar I should be somewhat more than that calculated from the artificial electron belt, because slow proton damage probably contributes

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somewhat to the degradation (private communication from W. Brown, Bell Telephone Laboratories).

Injun, Tiros V, and other satellites continue to function. Injun has a low duty cycle and Tiros V shows some solar cell degradation. Film badge dosimeter measurements have been made on several U. S. Department of Defense satellites. About 10 r/day was measured\* inside 1.5 gm/cm<sup>2</sup> of shielding. In order to compare this radiation dose with the predictions in Table 1, correction must be made for the shielding. To do this we perform a calculation like that done for the Traac GM counter to get *f*, the fraction of electrons that penetrate the wall. Values of *f* have been calculated for different thicknesses of shield by using the relationship  $R = 0.526 E - 0.094$  and the associated rough-straggling transmission curves in Figure 3. Figure 9 shows a plot of 1/*f* as a function of

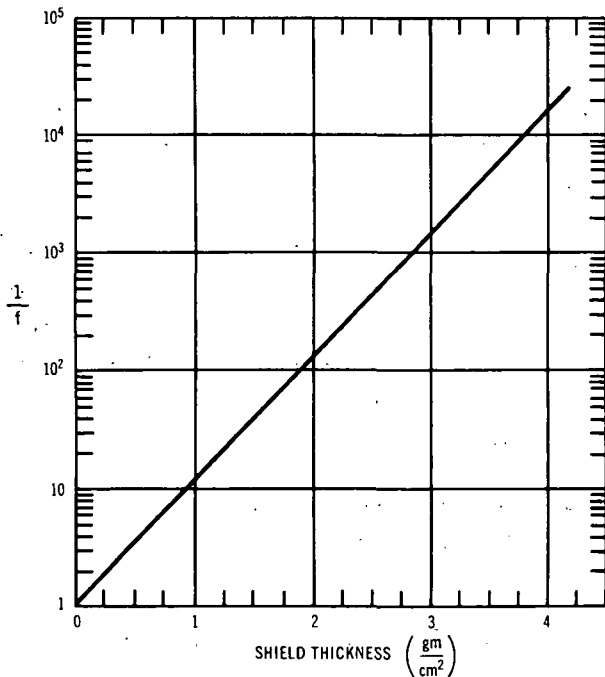


Figure 9—The fraction of fission electrons that penetrate different shield thicknesses.

\*This number is probably uncertain by a factor of 2 because of the nature of the radiation causing the film blackening (private communication from R. Moffet, Lockheed Missile and Space Company, Palo Alto, California). Also an unexplained up-down difference of a factor of 3 exists in the data.

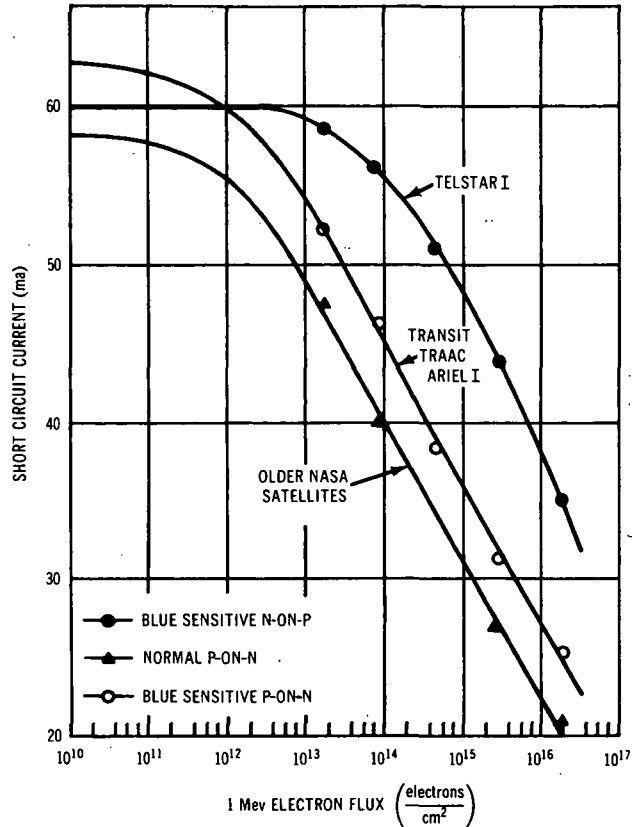


Figure 8—Solar cell damage curves.

shield thickness. This is really only true for Al but for lack of better information we will use it for other materials too. For 1.5 gm/cm<sup>2</sup> we get *f* = 1/50 for normal incidence particles. To correct for a distribution of incidence angles we will say roughly that about half as many get through (private communication from W. Bethe, Cornell University). Also, 2π steradians are covered by a much thicker shield so that the total factor *f'* = 1/200. This would mean that 10 r/day × 200 = 2000 r/day were incident on the outside of the vehicle. This agrees quite well with numerical calculations.

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## MANNED FLIGHT

For a Mercury capsule orbit with an apogee of 264 km the total flux encountered in six orbits would be  $0.24 \times 10^{10}$  electrons/cm<sup>2</sup> outside the vehicle (Table 4). If the apogee is lowered by 30 km (to 234 km) the total flux for 6 orbits is reduced to  $0.17 \times 10^{10}$  electrons/cm<sup>2</sup>. If the apogee is raised by 30 km (to 294 km) the total flux for 6 orbits is increased to  $0.45 \times 10^{10}$  electrons/cm<sup>2</sup>.

## PARTICLE TIME HISTORIES

One of the important problems to answer about the new belt is how long it will last. The currently intense regions will last a number of years, according to present indications. At low altitudes the fluxes have already decayed a lot. According to Ariel I and Traac data, outside the  $10^5$  contour of the B-L plot in Figure 6 the fluxes decayed several orders of magnitude in a few days.

Injun has noted some decay at 1000 km (private communication from B. O'Brien, State University of Iowa). At  $L = 1.18$  and  $B = 0.191$  there is a decay factor of about 2 from +10 to +1000 hr. For the same  $L$  and time interval for  $B = 0.206$  there is a decay factor of 4. Injun saw no marked change in flux as a result of a modest size magnetic storm.

The only decay process we understand well enough to calculate is coulomb scattering. Particle time histories have been calculated for coulomb scattering and characteristic times determined (Reference 5). The time to reach a scattering equilibrium (which is also about the time for this equilibrium to decay to  $1/e$  intensity) for different  $L$  values is listed in Table 5. Welch, Kauffman, et al. (Reference 5) first calculated these for solar maximum atmospheric densities and now, assuming that the density is less by a factor of 10, we get the values in Table 5.

Table 5  
Time Until Scattering Equilibrium for Different Values of  $L$ .

$L$	Calculated $\tau$ (days)	Measured $\tau$
1.20	10	~ 1 month
1.25	150	-
1.30	1500	-
1.35	~3000	-
1.40	~10,000	-

Table 4

Flux per Orbit for a Mercury Capsule at an Altitude of 264 km.

Orbit	Flux (electrons/cm <sup>2</sup> )
1	$5.0 \times 10^6$
2	$2.1 \times 10^7$
3	$4.8 \times 10^7$
4	$2.9 \times 10^8$
5	$6.4 \times 10^8$
6	$1.4 \times 10^9$

The densities are not well known and the calculated times may be wrong by a factor of 5 or more. The calculated variation with L, however, should be fairly good. Although the Injun data do not show the expected variation with L, they do show that the calculated times are of the right order of magnitude. The times show that the high flux region should last even through the next solar maximum if coulomb scattering is the principle loss process.

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Appendix A

### Shielding and Radiation Doses

Some crude calculations on shieldings and dosages are given here. Mr. William Gill of Marshall Space Flight Center is doing more complete and quantitative work on this subject and for better information he should be contacted. This appendix is included only for the sake of completeness.

One consideration that is important in some shielding calculations is bremsstrahlung. The doses delivered by the x rays made by bremsstrahlung will be larger than the direct electron doses for large shield thicknesses.

The fraction of the energy of an electron that goes into bremsstrahlung may be calculated from:\*

$$\frac{E_{\text{brem}}}{E_{\text{ion}}} = \frac{ZE^2}{1600}$$

where Z is the atomic number of the material involved. For the fission energy spectrum the average energy is about 1 Mev;

$$\frac{E_{\text{brem}}}{E_{\text{ion}}} = \begin{array}{cccc} \text{C} & \text{Al} & \text{Fe} & \text{Pb} \\ 0.004 & 0.008 & 0.015 & 0.050 \end{array}$$

The energy spectrum of the x rays will be something like that in Figure A1. There will be a very few x rays up to 8 Mev, but not many over 2 or 3 Mev. The low energy x rays (below about 100 kev) will be absorbed in the shielding. This will remove about half the total energy in the x rays. The resultant transmitted energy spectra will have a peak at about 1/2 Mev (Figure A1). The x rays transmitted through the shield will be quite penetrating. Their mean free path will be roughly 20 gms/cm<sup>2</sup>. This means two things. First, they will be hard to absorb, and therefore it will take a lot more shielding to absorb them. Second, because they are hard to absorb, they will not be counted efficiently by a particle counter and also will result in less radiation dose.

We can now calculate crudely the counting efficiency of the Injun (1961 02) 213 GM counter. From Figure 9 of the body of this report we see

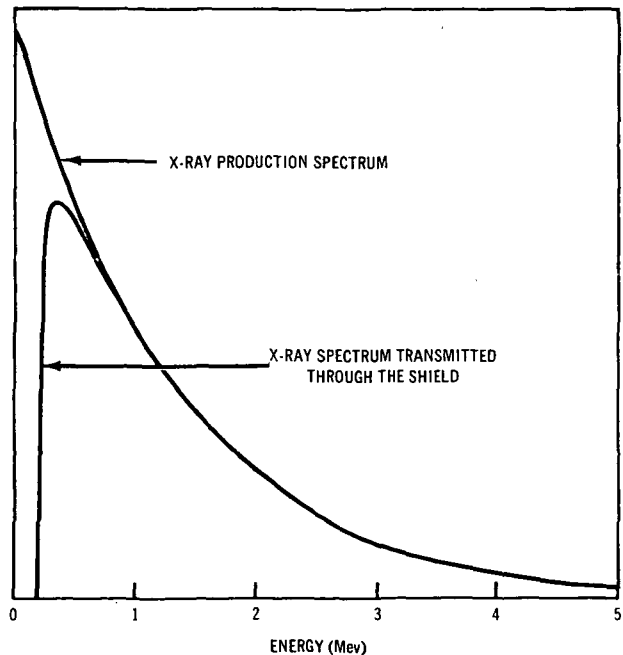


Figure A1-A crude bremsstrahlung x-ray energy spectrum.

\*Fermi, E., "Nuclear Physics," A course given by Enrico Fermi at the University of Chicago, Notes compiled by Orear, J., Rosenfeld, A.H., and Schluter, R.A., University of Chicago Press, revised edition, 1950.



that it would only count about 1/20,000 of the fission electrons directly. But we find that 0.05 of the energy is converted to bremsstrahlung, of which half is absorbed in the shield. The mean energy of these x rays will be about 1/2 Mev. A normal GM counter will detect these x rays with about 1 percent efficiency. This gives

$$(0.05)(1/2)(0.01) = 1/4000 ,$$

for the fraction of the electrons counted via bremsstrahlung. The agreement with the actually determined factor of 4000 is fortuitous here. This calculation is not extremely accurate but it does show that the Injun counter counts electrons via bremsstrahlung with about the observed efficiency.

### Manned Flight

The effects of the new radiation belt on manned flights must be considered. For the Mercury project the total flux that would be encountered for a six orbit mission with the MA-7 (1962  $\tau$ 1) orbit is  $0.24 \times 10^{10}$  electrons/cm<sup>2</sup> outside the vehicle, or 100 r ( $2.5 \times 10^7$  electrons/cm<sup>2</sup> = 1 r). The shielding of the vehicle is given in Table A1 (private communication from Carlos Warren, Marshall Space Flight Center).

Table A1  
Shielding and f Values on a Spacecraft, for Different Solid Angles.

Solid Angle (ster)	Shield Thickness (gm/cm <sup>2</sup> )	f
0.87	0.58	1/4.4
0.53	0.85	1/8.0
0.35	1.27	1/21
0.07	2.57	1/400
rest	> 5	0

The total shielding factor is:

$$\left(\frac{0.87}{4\pi}\right) \frac{1}{4.4} + \left(\frac{0.53}{4\pi}\right) \frac{1}{8.0} + \left(\frac{0.35}{4\pi}\right) \frac{1}{21} + \left(\frac{0.07}{4\pi}\right) \frac{1}{400} = 0.022 .$$

This gives a dose inside the capsule of  $100 \times 0.022 = 2.2$  r for the six orbit mission. Actually the dose will be less than this, about 1.2 r, because of the nonperpendicular incidence of the electrons. This is the best current estimate of the dose the astronaut will take. This will be a skin dose and will not penetrate very far — only a few centimeters. Also, part of the body is protected more than this, so the 1.2 r is not true for the whole body.

Almost all of the dose would be received in the South Atlantic "hot spot" (see Figure 8 of the body of the report) and would occur mainly on orbits 4, 5, and 6. The breakdown of the 1.2 r dose inside the capsule by orbits is given in Table A2.

The calculation which has been made to give the dose inside the Mercury capsule is probably correct to a factor of 3. The fluxes in the South Atlantic "hot spot" (see Figure 7 in the body of the report) are probably correct to a factor of 2. We will get more information from Orbiting Solar Observatory I (1962 ζ 1) and a future Department of Defense flight to help confirm this. The shielding calculations should be checked, but they are also probably correct to a factor of 2. Fortunately, the calculated dose is comfortably below the mission tolerance.

Table A2  
Radiation Dose per Orbit for a  
6 Orbit Flight on the MA-7 Orbit.

Orbit	Dose (r)
1	0.003
2	0.014
3	0.028
4	0.14
5	0.32
6	0.69

For the Apollo project the problem is different. Apollo will go out through the most intense part of the artificial belt in a period of about 1000 seconds. The approximate dose outside the Apollo capsule will be  $2000 \text{ sec} \times 10^9 \text{ electrons/cm}^2\text{-sec} = 2 \times 10^{12} \text{ electrons/cm}^2 = 80,000 \text{ r}$  due to the artificial belt for the lunar round trip. The shielding is quite thick so that the major dose will be received from bremsstrahlung x rays. Assuming most of the wall thickness to be Fe we get  $0.015 \times 1/2$  of the energy through the wall, but only through  $2\pi$  steradians, because the rest will be covered with a very thick absorber. The x rays that will get inside will be of an energy of  $\sim 1/2 \text{ Mev}$  and have a mean free path of about  $20 \text{ gm/cm}^2$ . To calculate the radiation dose, we want the energy deposition per gram so we divide the energy in the x-ray beam by 20 to get the fraction absorbed per gram. This gives:

$$(80,000 \text{ r}) \times (0.015 \times 0.5) \times (0.5) \times (1/20) = 16 \text{ r (whole body dose inside Apollo capsule).}$$

calculated external dose	fraction of energy through Apollo wall as x rays	fraction of solid angle	fraction of energy deposited per gram
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This does not seem serious, but it is a crude calculation and should be performed to a greater degree of accuracy. The dose could be cut by 3 by using about 1 inch of  $\text{CH}_2$  on the outside of the vehicle to cut the x-ray yield.